



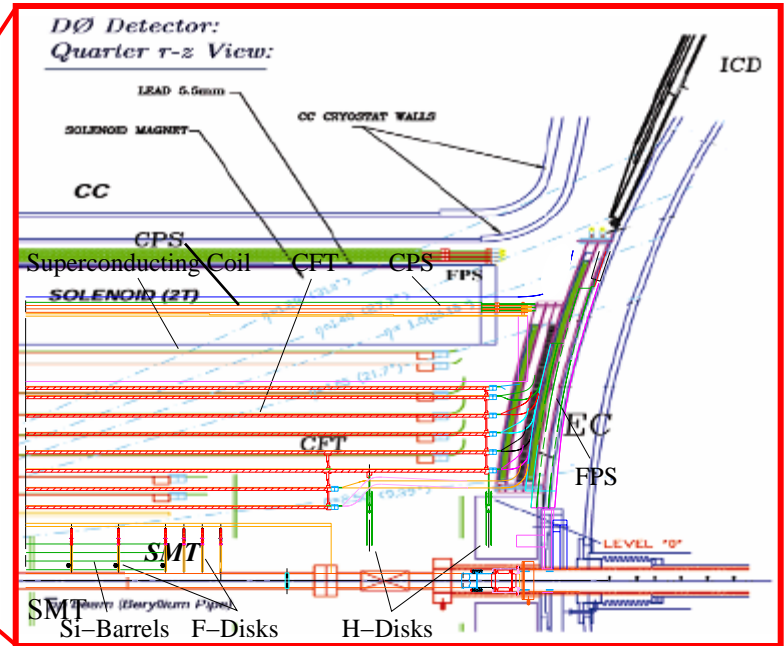
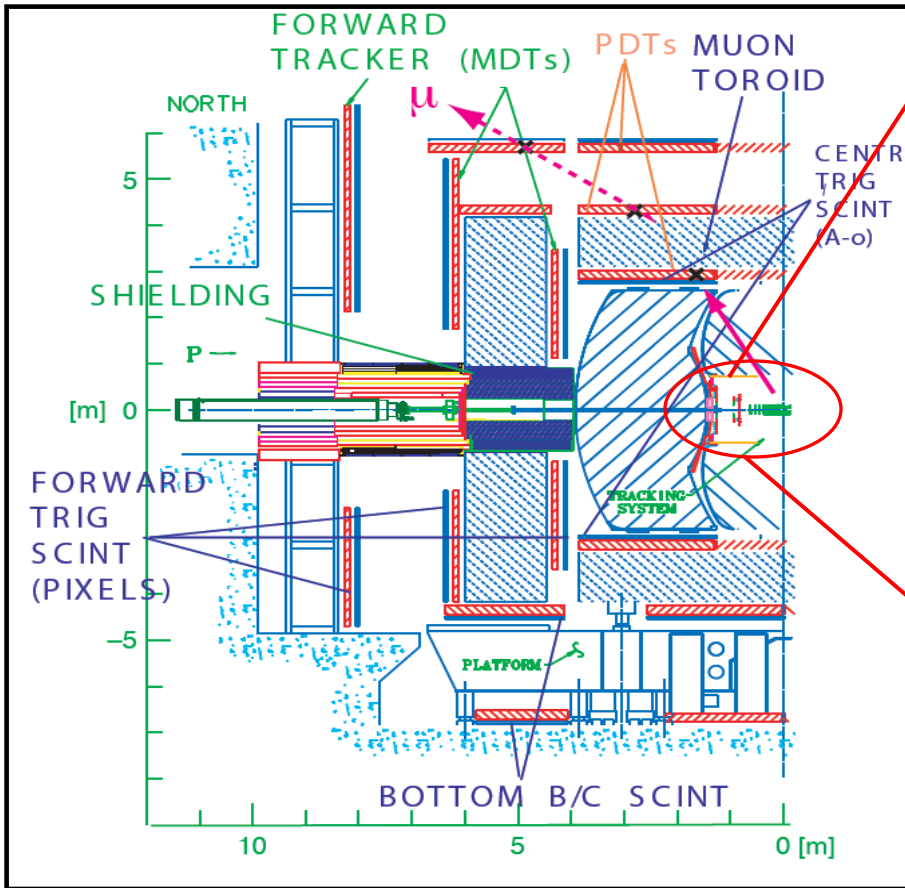
# B Hadron Properties at DØ

H. Eugene Fisk for the DØ Collaboration

- Introduction
- Orbitally Excited b-mesons
- Bc Mass and Lifetime
- Cascade B
- Conclusions



# DO Detector



- General purpose detector
- Excellent coverage of Tracking and Muon Systems ( $|\eta| < 2$ )
- Excellent vertex resolution
- 2T Solenoid



# Why B spectroscopy?

- b-quarks => heavy bound state mesons and baryons. The mesons serve as the "hydrogen atoms" of the quark world and the particle theorists love them because they can calculate!
- b-states measurements provide sensitive tests of potential models, heavy quark effective theory (HQET) and even tests of lattice gauge theory. Complementary to charm states.
- B-quark production cross-sections are large :  $\sim 1 \mu\text{b}$  for  $P_t > 5 \text{ GeV}/c$ , can provide  $\sim 1200$  events/day at  $L \sim 100 \text{ E30}$ .
- Tevatron energy provides access to b-states that are not accessible at e+e- b-Factories.

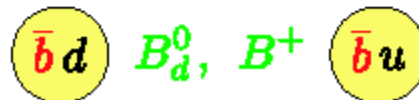
Outline:

Heavy Mesons

Heavy Baryons



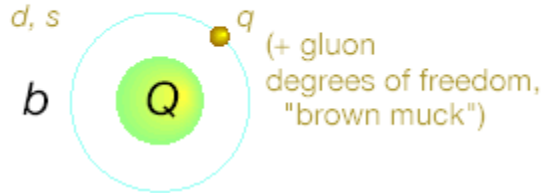
→ Complementary to  $\Upsilon(4S)$  B factories





# Orbitally-excited B mesons

Hydrogen atom of strongly interacting systems

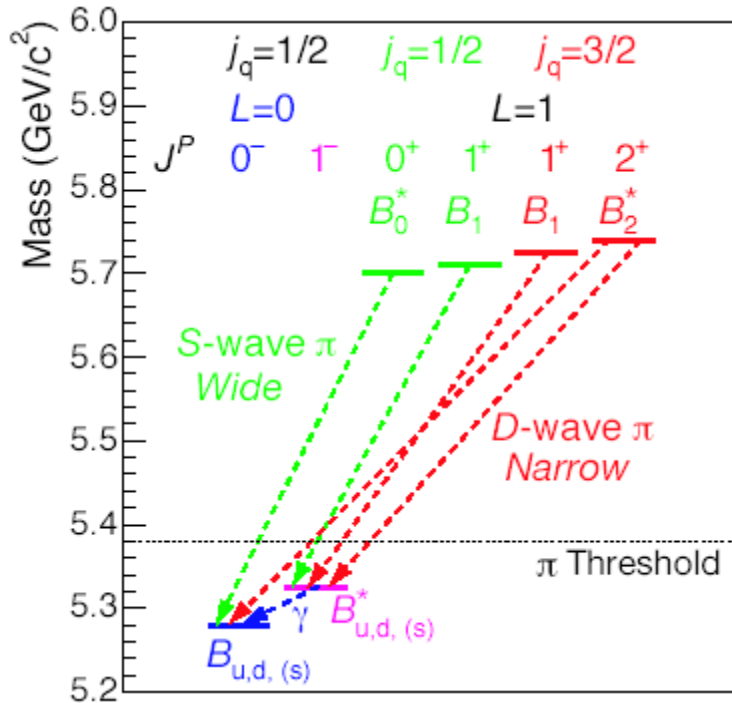


$$\vec{J}_q = \vec{s}_q + \vec{L} \quad \vec{J} = \vec{s}_q + \vec{J}_q \quad L = 1$$

$$j_q = 1/2 \quad J = 0, 1 \quad \boxed{B_0^*, B_1^*} \quad \boxed{B_{s0}^*, B_{s1}^*}$$

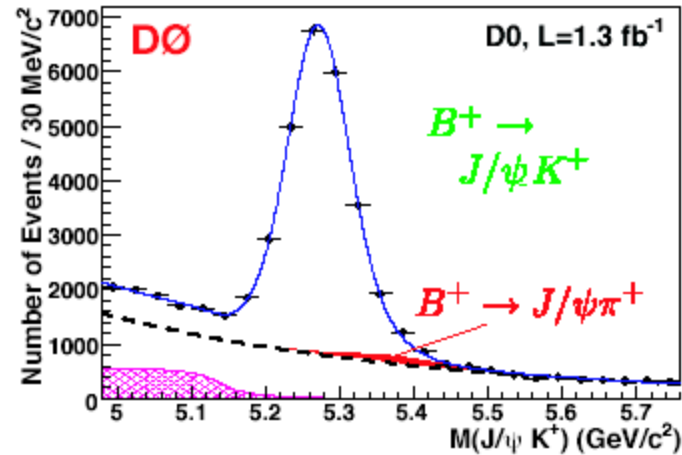
$$j_q = 3/2 \quad J = 1, 2 \quad \boxed{B_1, B_2^*} \quad \boxed{B_{s1}, B_{s2}^*}$$

Collectively referred to as:  $B^{**}$   $B_s^{**}$   
 $B_J$   $B_{sJ}$



$B_1$  or  $B_2^* = B^{(*)}\pi^- \Rightarrow J/\psi K^+ \pi^-$   
 $J/\psi \Rightarrow \mu^+ \mu^-$  and  $B^* \Rightarrow B \gamma$

45.8 MeV  $\gamma$  not observed

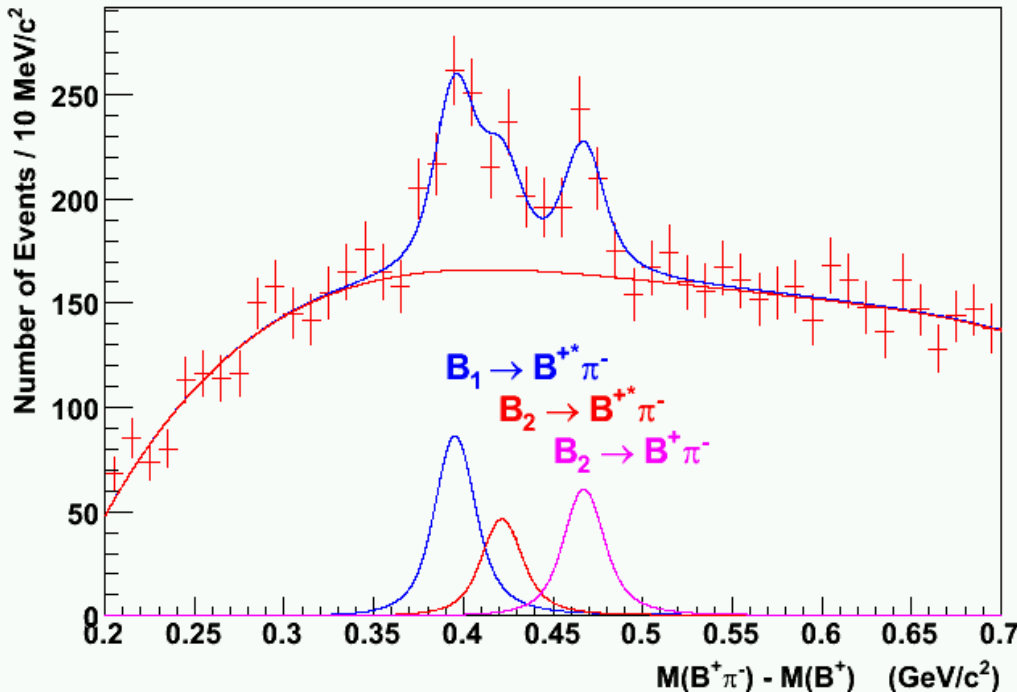




# B<sup>\*\*</sup> Mesons



D0, 1.3 fb<sup>-1</sup>



$$M(B_1) - M(B^+) = 441.5 \pm 2.4 \pm 1.3 \text{ MeV}/c^2$$

$$M(B_2^*) - M(B_1) = 26.2 \pm 3.1 \pm 0.9 \text{ MeV}/c^2$$

$$M(B_1) = 5720.6 \pm 2.4 \pm 1.4 \text{ MeV}/c^2$$

$$M(B_2^*) = 5746.8 \pm 2.4 \pm 1.7 \text{ MeV}/c^2$$

$$N(B^{**}) = 662 \pm 91 \pm 140 \text{ events}$$

$$\Delta_1 = M(B_1) - M(B^*),$$

$$\Delta_2 = M(B_2^*) - M(B^*),$$

$$\Delta_3 = M(B_2^*) - M(B) = \Delta_2 + 45.8 \text{ MeV}$$



## B\*\* Mass Fitting

$$F(\Delta M) = F_{\text{sig}}(\Delta M) + F_{\text{bckg}}(\Delta M)$$

$$F_{\text{sig}}(\Delta M) = N\{f_1 D(\Delta M, \Delta_1, \Gamma_1) \\ + (1 - f_1)[f_2 D(\Delta M, \Delta_2, \Gamma_2) \\ + (1 - f_2) D(\Delta M, \Delta_3, \Gamma_2)]\}$$

$F_{\text{back}}(\Delta M) = 4^{\text{th}}$  order polynomial.  $D(x, x_0, \Gamma)$  is the convol'n of a relativistic Breit-Wigner with Gaussian.

$$M(B_1) = 5720.6 \pm 2.4 \pm 1.4 \text{ MeV}/c^2$$

$$M(B_2^*) = 5746.8 \pm 2.4 \pm 1.7 \text{ MeV}/c^2$$

$\chi^2/\text{d.o.f.} = 33/40$ ; w/o  $B_J$  signal  $\chi^2/\text{d.o.f.} = 97/45$ .

Peak structure significance  $> 7\sigma$ .

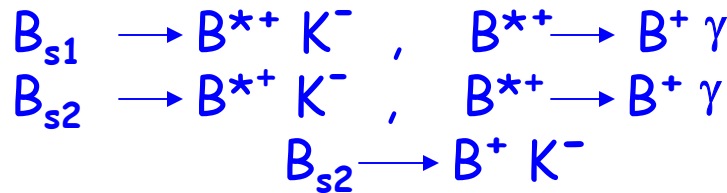
Fitting with only one peak,  $\chi^2/\text{d.o.f.} = 54/42$  or  $4\sigma$ .

With  $B_2^* \Rightarrow B^*\pi$  removed,  $\chi^2/\text{d.o.f.} = 41/41$ . Three peaks preferred, which agrees with theory, but statistically marginal.



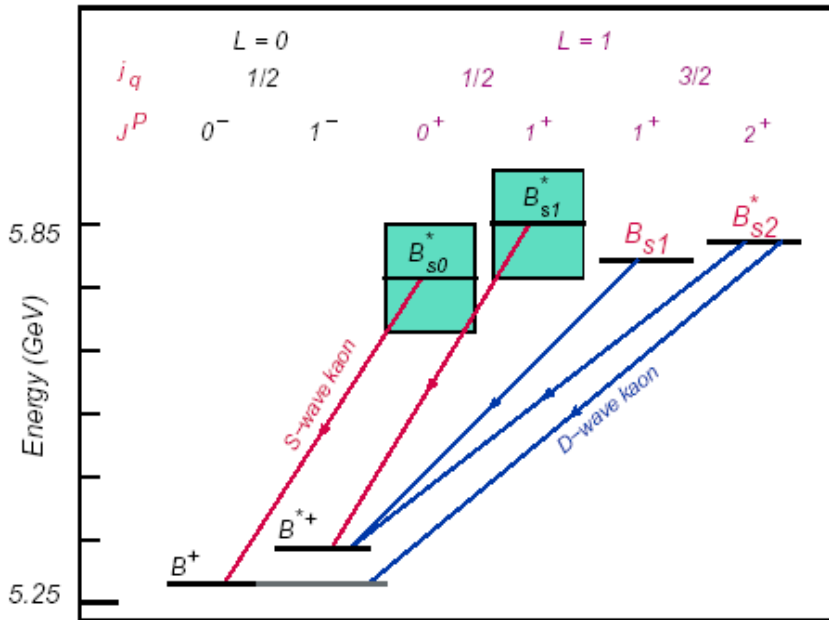
# Obs. of the Excited $B_{s2}$ Meson ( $b\bar{s}$ )

Use the same events just discussed in the  $B_1$  and  $B_2^*$  analysis, with the assignment of the charged track as a K meson. Look for:

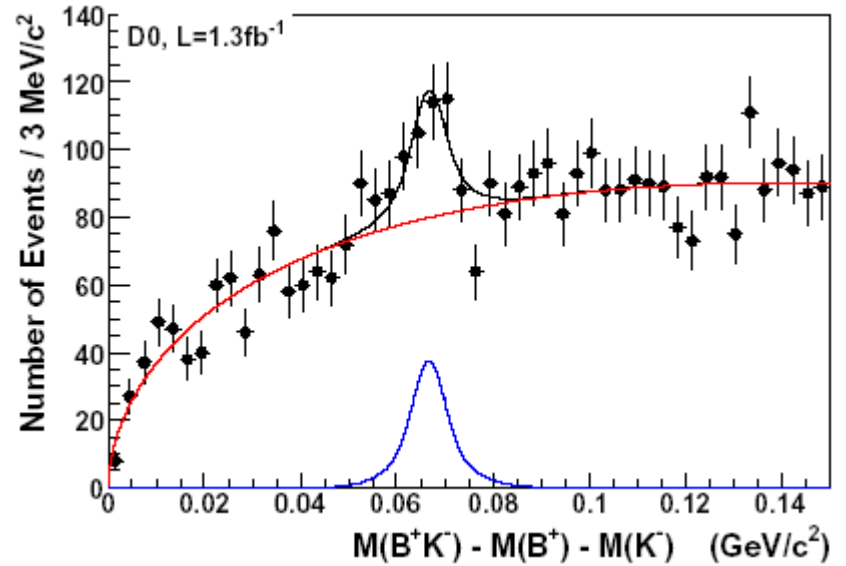


20,915 +/- 293 +/- 200  $B^+$  candidates  
 Additional track  $p_{\perp} > 0.6 \text{ GeV}/c$  w/charge opposite to the  $B$  charge.

Plot:  $\Delta M = M(B^+K^-) - M(B^+) - M(K^-)$



PRL 100,082002 (2008)





# $B_{s1}, B_{s2}^*$ Search Results

Of the three decays through which the  $B_s J$  states can reach the ground state  $B^+$ , one or more may be forbidden if the excited state mass is smaller than the mass of the decay products. One obvious bump at  $\Delta M = 67 \text{ MeV}/c^2$ . This is the highest energy transition:  $B_{s2} \rightarrow B^+ K^-$ . Because this mass is close to threshold its width  $\Gamma$  is expected to be  $\sim 1 \text{ MeV}/c^2$ . The apparent width is due to our mass resolution, so  $\Gamma$  is fixed at  $1 \text{ MeV}/c^2$ . Fit to a single B-W with Gaussian expt'l resolution:

## Single Peak Fit:

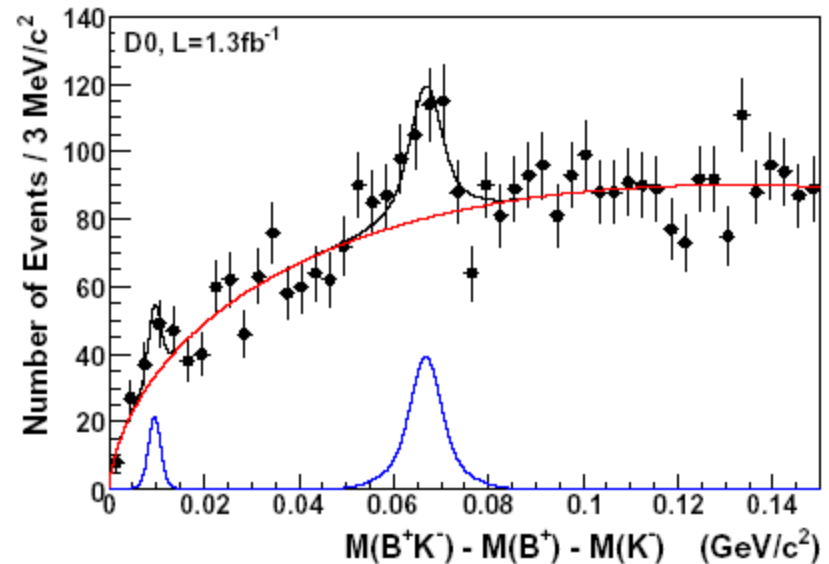
$$\begin{aligned}\Delta_0 &= M(B_{s2}^*) - M(B^+) - M(K^-) \\ &= 66.7 \pm 1.1(\text{stat}) \text{ MeV}/c^2.\end{aligned}$$

w/o  $B_{s2}$  signal fit log L shows  $4.8\sigma$ .

## Double Peak Fit:

$\Delta M(B_{s1}) = 11.5 \pm 1.4(\text{stat}) \text{ MeV}/c^2$ .  
w/o  $B_{s1}$  sig, log L shows a  $2.7\sigma$  decrease so it is inconclusive.

## Double Peak Fit

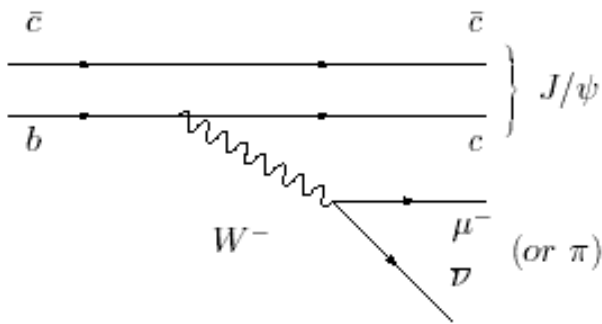






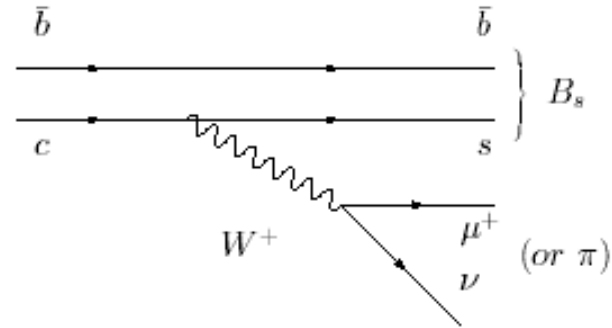
# $B_c$ Meson ( $b, \bar{c}$ ): Mass and Lifetime

- Two heavy quarks; decays via  $b$  or  $c$  to  $W$ ; also annihilation.



$c$  as spectator (this talk)

$J/\psi \mu \nu$  or  $J/\psi \pi$



$b$  as spectator (not seen)

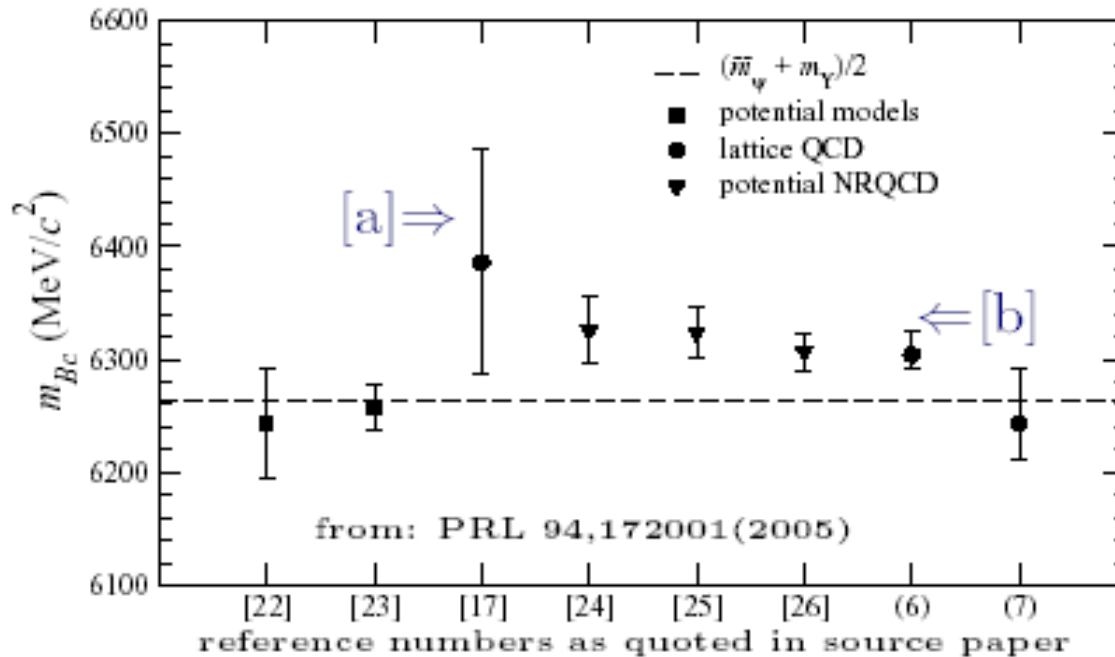
$B_s \pi$  or  $B_s \mu \nu$

Expect the lifetime to be  $\sim 1/3$  of other  $B$ -mesons.

For mass measurement  $J/\psi \pi$  final state has no missing particles.  
Has di-muon plus charged pion vertex. Can tune up on a state with similar topology:  $B^+ \rightarrow J/\psi K^+$ .



# $B_c$ Mass Predictions



[a] Omitting sea quarks, (quenched approx.).

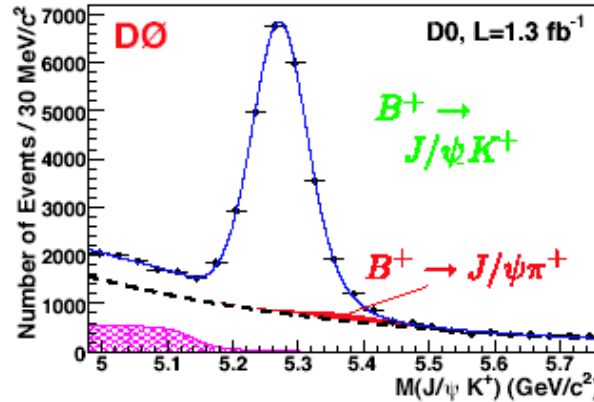
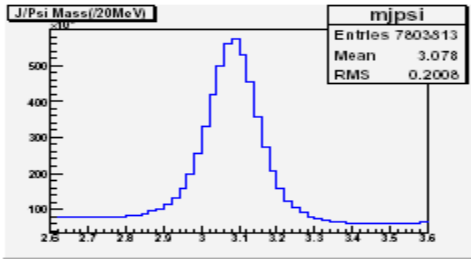
[b] Add 2+1 sea flavors,  $u, d$  as light as possible, and strangeness mass.

[a] H.P. Shanahan, P. Boyle, C.H.T. Davies, H. Newton  
PLB 453, 289 (1999) 6386 +/- 9 +/- 15 +/- 98 MeV/c<sup>2</sup>

[b] I.F. Allison, C.T.H. Davies, A. Gray, A.S. Kronfeld, P. B. Mackenzie, J.N. Simone  
PRL 94, 172001 (2005) 6304 +/- 4 +/- 11 (+18/-0) MeV/c<sup>2</sup>



# Data Cuts: Learn on $B^+ \rightarrow J/\psi K^+$ signal and sidebands

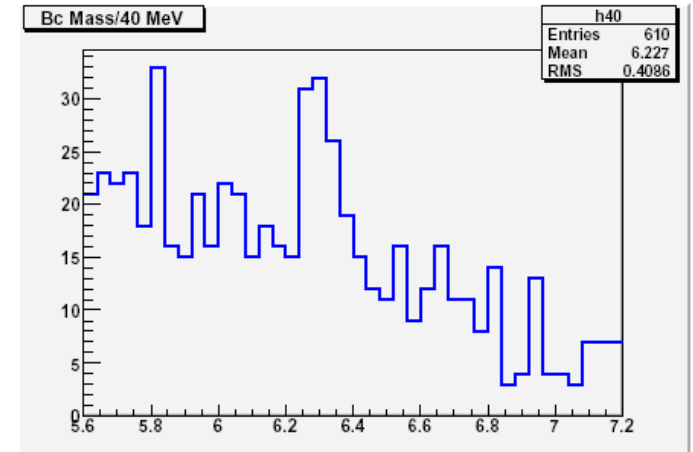
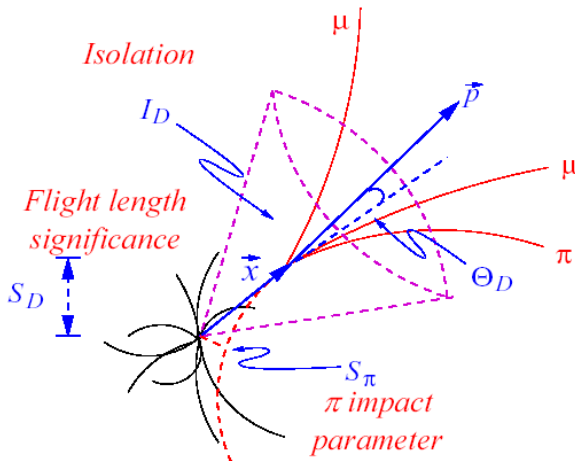


## Pre-selection Cuts:

- $P_T J/\psi > 4 \text{ GeV}/c$
- $P_T 3^{\text{rd}} \text{ particle} > 1.5 \text{ GeV}/c$
- $P_T(B) > 5 \text{ GeV}/c$

$J/\psi \rightarrow \mu^+\mu^-$   
7.6M events

$(J/\psi K^+)$  Mass  
37K events



$J/\psi \pi$  Mass

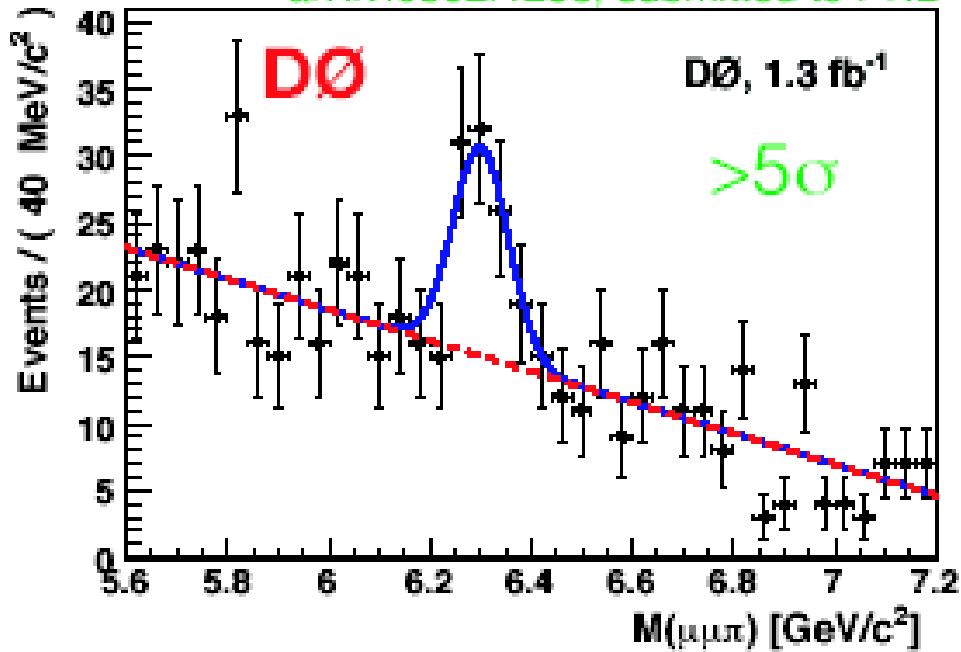
## Data based cuts:

- IPsigB: 3-D Impact Par Signif of B wrt PV;
- IPsig3: 3-D Impact parameter - third track;
- $\chi^2$ : formation of the 3 particle vertex;
- Sxy: 2-D Decay length Significance;
- Cxy: 2-D point cos for the B cand w.r.t. PV



# $B_c$ Mass

accepted by  
arXiv:0802.4258, submitted to PRL



$$m(B_c^+) = 6300 \pm 14 \pm 5 \text{ MeV}$$



# $B_c \Rightarrow J/\psi \mu + X$ Lifetime Measurement I

arXiv:0805.2614v1 [hep-ex]

- Integrated Luminosity of  $\sim 1.3 \text{ fb}^{-1}$ .
- Decay length from primary vertex to  $J/\psi \mu$  vertex.
- The  $B_c$  signal is established by using mass fits following decay length requirements.
- MC samples of signal and backgrounds are used to model mass and decay length distributions.
- $J/\psi$  sample from reconstructed opposite sign  $\mu$ 's w/pt  $> 1.5 \text{ GeV}/c$  and di-muon mass in the range  $2.90$  to  $3.26 \text{ GeV}/c^2$
- Add a third muon whose track forms a good vertex with the  $J/\psi$ .
- Find 14,753 events with  $J/\psi \mu$  masses between 3 and  $10 \text{ GeV}/c^2$ .

The  $J/\psi \mu$  mass can be used to characterize the components that make up the sample:

- (1) Signal - SI
- (2) Real  $J/\psi$  with a "fake" muon due to a track - JT (small component)
- (3) Fake  $J/\psi$  from combinatorial background - CB
- (4) Real muon forms a vertex with a real  $J/\psi$ , neither from a  $B_c$  - JM
- (5)  $B^+ \rightarrow J/\psi K^+$  with  $K \Rightarrow \mu \nu$  in flight - DK
- (6)  $c$  pair where a prompt  $J/\psi$  is associated with a muon - PR

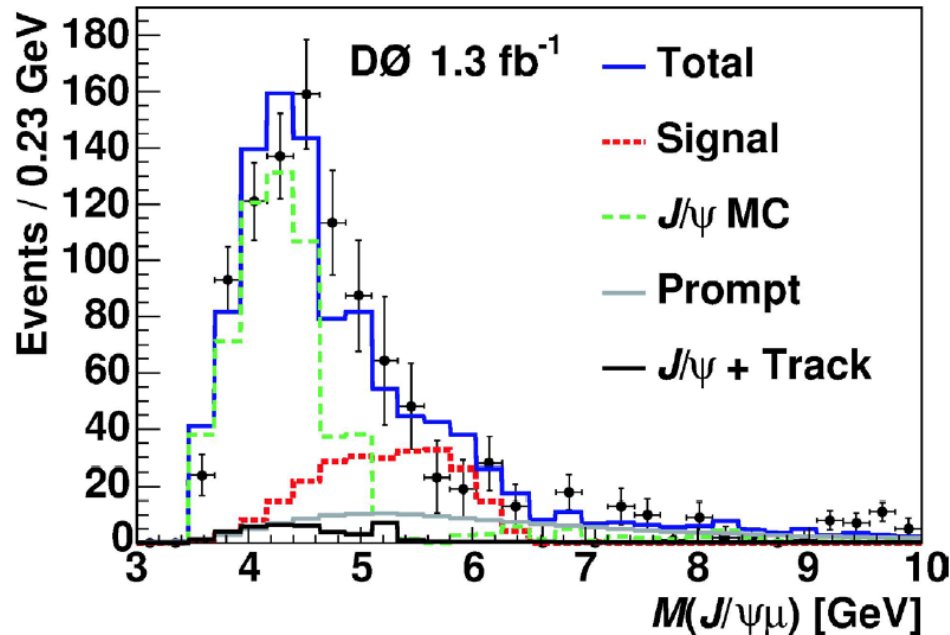
The  $J/\psi \mu$  invariant mass is used to model the six components. And the decay length is also a parameter. The lifetime of the  $B_c$  is not measured because of the missing  $\nu$ .



## $B_c \Rightarrow J/\psi \mu$ Lifetime Measurement (2)

Use the Visible-Proper Decay Length (VPDL) =  $L_{xy} \frac{m(B_c^\pm)}{p_T(J/\psi \mu)} = \frac{c\tau}{K}$

to define the K factor distribution that is then applied statistically when extracting the  $c\tau(B_c)$ .

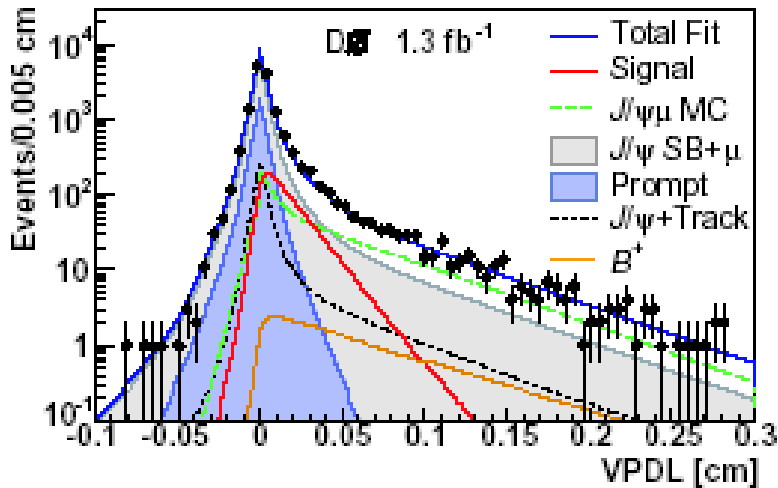


Fit to the  $J/\psi + \mu$  vertex with the  $J/\psi$  mass sideband and  $B^+$  components subtracted and decay length significance  $L_{xy}/\sigma(L_{xy}) > 4$  required.



# $B_c \Rightarrow J/\psi \mu$ Lifetime Measurement (3)

- Both the  $M(J/\psi \mu)$  and VPDL distributions can be fitted simultaneously after the various components of backgrounds are established.



Summary of estimated systematics	
Systematic Source	$\Delta\tau$ (ps)
Mass model uncertainty	$\pm 0.021$
Lifetime model uncertainty	$\pm 0.022$
Signal feed-down fraction	$\pm 0.005$
Alignment	$\pm 0.006$
Total	$\pm 0.032$

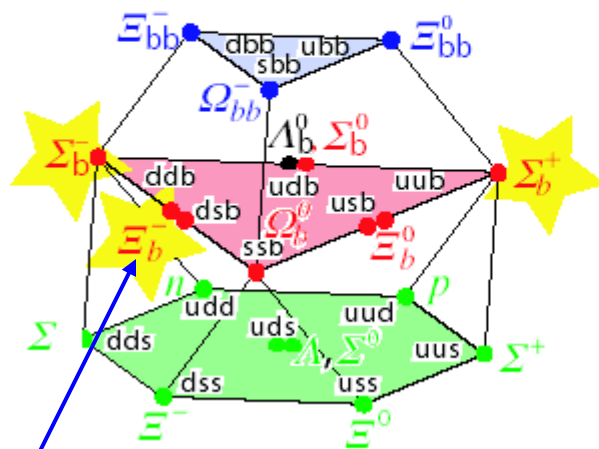
**World's best!**

$$\tau(B_c) = 0.448 +0.038/-0.036(\text{stat}) \pm 0.032(\text{sys}) \text{ ps}$$

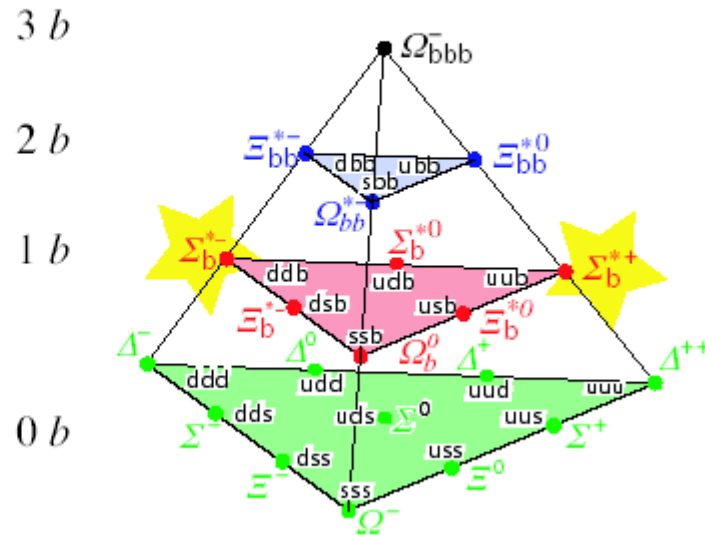


# b Baryons: Where are they?

$J=1/2$   $b$  Baryons



$J=3/2$   $b$  Baryons

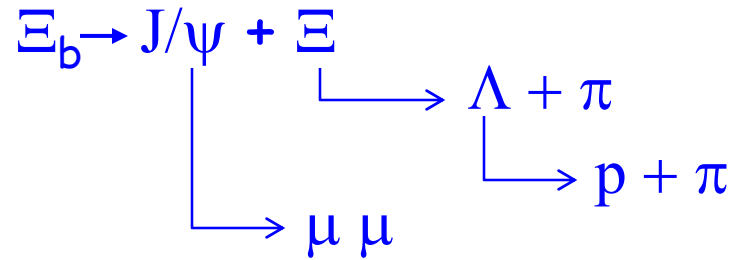
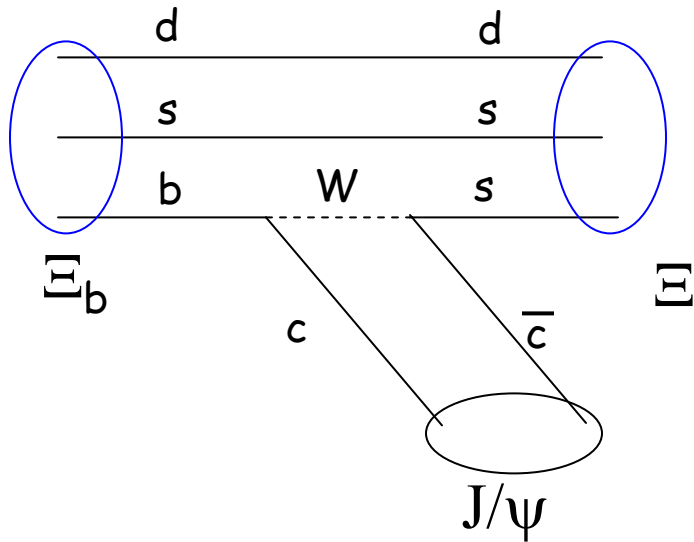


Search for the Cascade  $b$

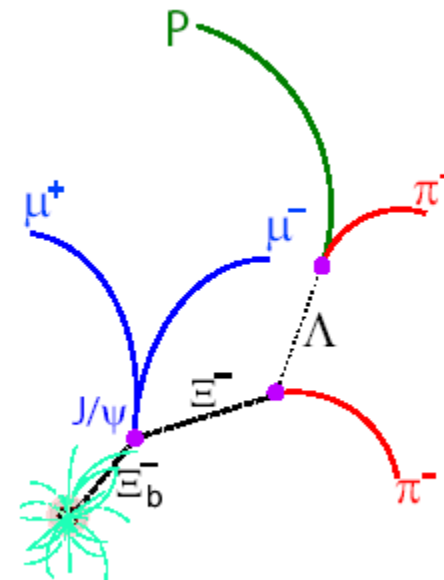




# Cascade b: (d, s, b) or $\Xi_b$



Challenging to connect all the dots!  
 Start with di-muons, then developed  
 downstream search for  $\Lambda$  search.  
 Kinematics constraints were then applied.





# $\Xi_b$ Candidate Selection

J/ $\psi$  requirements:

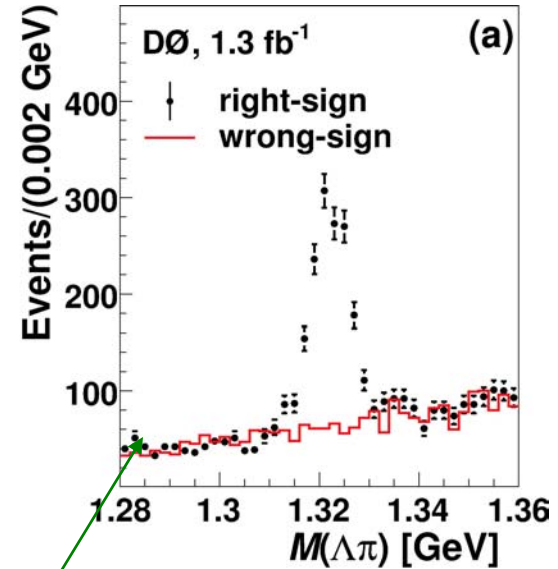
- Opposite sign  $\mu$ 's w/common origin;
- Matching  $\mu$  track in CT and  $\mu$  PDTs;
- $P_t(\mu) > 1.5$  GeV/c;
- $2.5 < M(\mu, \mu) < 3.6$  GeV/c;
- If more than one cand., take closest.

Other requirements:

- Primary vertex;
- Reprocess evts using improved eff. for finding downstream vertices/tracks.
- For  $\Lambda$  candidates assume the high  $p_t$  track is a proton (MC says  $\sim 100\%$  OK)
- $1.105 < M(p, \pi^-) < 1.125$  GeV/c;

$\Xi^-$  and additional  $\pi^-$  requirements:

- $\Lambda$  and  $\pi$  common upst. vertex; IP reqs
- Same charge sign for  $\pi_\Lambda$  and the other  $\pi$ ;
- Wrong sign  $\pi_\Lambda$  and other  $\pi \rightarrow$  bckgd level.

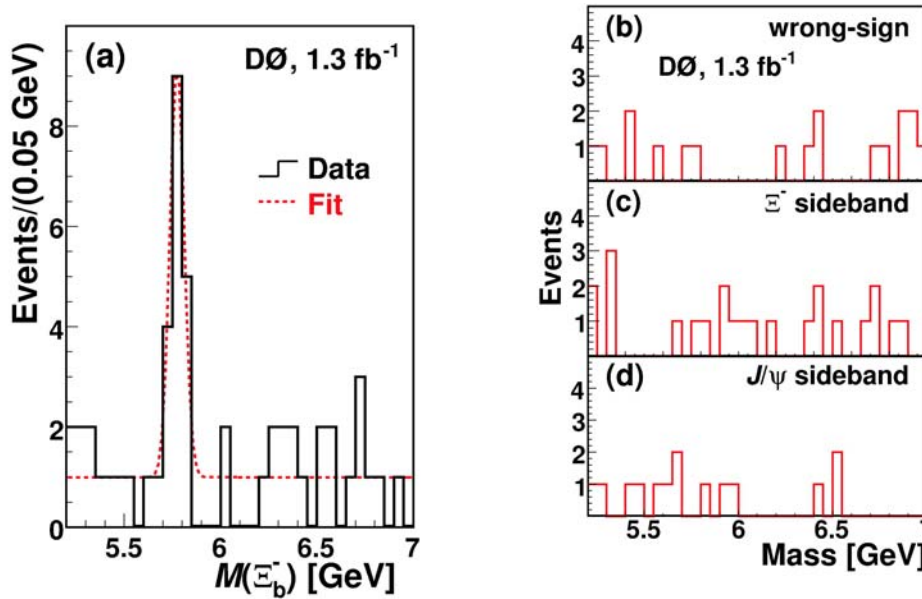


$\Xi_b \rightarrow J/\psi \Xi$  Conditions:

- For  $p(\Lambda)$ :  $p_t > 0.7$  GeV/c
- B decays suppressed  $l_{\Xi} > 5$ mm (avg  $l_{\Xi}$  in MC is 4.8cm)
- $\cos(\theta_{col}) > 0.99$  (MC studies) (Reduce Bkgd 56%; -1.7% Sig)
- $l_{\Xi_b}$  DK length sig  $> 2$  ( $l_{\Xi_b}$  MC)



# $M(\Xi_b)$ Mass: 1<sup>st</sup> Observation



$M(J/\psi \Lambda \pi^+)$  Wrong sign  $\pi$

$M(J/\psi \Lambda \pi^-)$  in  $\Xi^-$  sidebands  
 i.e.  $1.28 < M(\Lambda\pi^-) < 1.36 \text{ GeV}/c^2$  but  
 excluding  $1.305 < M_{\Xi^-} < 1.340 \text{ GeV}/c^2$

$M(J/\psi \Lambda \pi^-)$  in  $J/\psi$  Lower Sideband  
 Avoid upper sideband because of  $\psi'$ .  
 $2.5 < M(\mu^+\mu^-) < 2.7 \text{ GeV}/c^2$

$$M(\Xi_b) = M(J/\psi \Xi^-) - M(\mu^+\mu^-) - M(\Lambda\pi^-) + M_{\text{PDG}}(J/\psi) + M_{\text{PDG}}(\Xi^-)$$

Unbinned Likelihood fit gives:

$$= 5.774 \pm 0.011(\text{stat}) \pm 0.015(\text{syst}) \text{ GeV}/c^2$$

19 Candidate Events fitted with a Gaussian + flat background yields a signal/bkgd of 15.2/4.4 events.  
 Log likelihood ratio  $\sqrt{2\ln(L_{(s+b)}/L_b)}$   
 $\rightarrow 5.5\sigma$ , thus the probability of the signal being a bckg fluctuation is  $3.3 \times 10^{-8}$ ,  
 $\Rightarrow 2.2 \times 10^{-7}$  with all errors included.

Systematic Errors:

Event selection – Using multivariate technique using different variables for event selection leads to a conservative  $\pm 0.015 \text{ GeV}/c^2$  syst in the  $\Xi_b$  mass.

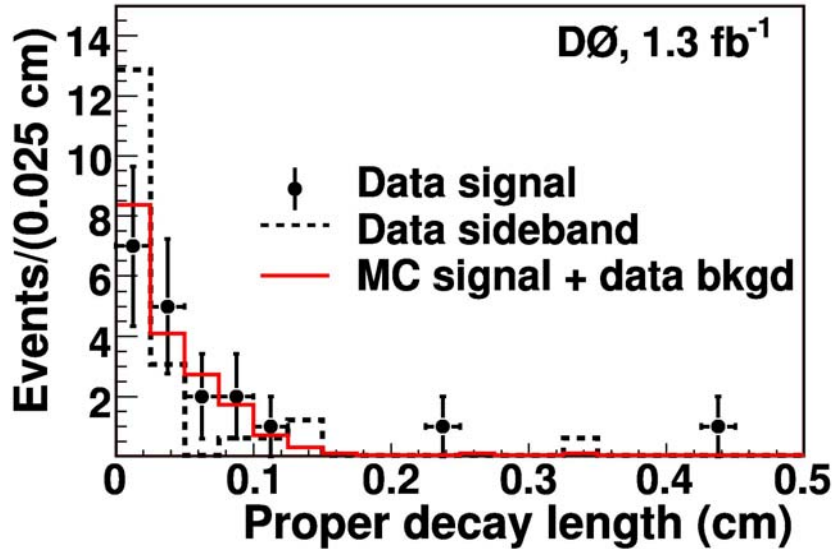
Signal & Background Models – Using the world average masses (PDG) shows a weak dependence on the track momentum scale. Verified with

$\Lambda_b \rightarrow J/\psi \Lambda$  and  $B^0 \rightarrow J/\psi K_s$  implies a mass uncertainty of  $\pm 0.002 \text{ GeV}/c^2$ .

CDF's  $M(\Xi_b)$  measurement is:  
 $5.792 \pm 0.0025 \pm 0.0017 \text{ GeV}/c^2$



# $\Xi_b$ Proper Decay Length Distribution



19 candidate events in the  $M_{\Xi_b} \pm 2.5\sigma$  window and sideband events, defined to be  $5\sigma$  away from the fitted mass, scaled to the number of events in the signal mass window, are included.

The expected distribution for 14.8 MC  $\Xi_b$  events and 3.6 background events are also shown.

K-S tests indicate that the distribution of signal events is favored over that of the sideband events, when compared to the MC generated  $\Xi_b$  events, by 5 to 1 odds.



## $\Xi_b$ ( $\sigma \times B$ ) Relative to $\Lambda_b$

$$\frac{\sigma(\Xi_b^-) \times B(\Xi_b^- \rightarrow J/\psi \Xi^-)}{\sigma(\Lambda_b) \times B(\Lambda_b \rightarrow J/\psi \Lambda)} = \frac{\varepsilon(\Lambda_b \rightarrow J/\psi \Lambda)}{\varepsilon(\Xi_b^- \rightarrow J/\psi \Xi^-)} \frac{N_{\Xi_b^-}}{N_{\Lambda_b}}$$

$$N(\Xi_b) = 15.2 \pm 4.4 +1.9/-0.4$$

$$N(\Lambda_b) = 240 \pm 30 \pm 12$$

$$\text{Efficiency ratio (MC)} = 4.4 \pm 1.3$$

$$\frac{\sigma(\Xi_b^-) \times B(\Xi_b^- \rightarrow J/\psi \Xi^-)}{\sigma(\Lambda_b) \times B(\Lambda_b \rightarrow J/\psi \Lambda)} = 0.28 \pm 0.09(\text{stat}) +0.09/-0.08$$



## Summary/Conclusions

B hadron properties are being understood with data from the Tevatron on many fronts:

- New states that excite us: e.g.  $\Xi_b$
- Orbitally excited mesons; we are making theorists happy.
- b-state lifetimes for previously unmeasured states,  $B_c$ ;
- We're testing lots of b-physics predictions like lattice gauge.
- The Fermilab collider is supplying much data for us. That's great! Our thanks to all who have done so well.